



University College Dublin
An Coláiste Ollscoile, Baile Átha Cliath

SEMESTER I EXAMINATION 2010/2011

ACM 40460
Dynamic Meteorology

Extern examiner: Professor Peter A Clark
Head of School: Professor Mícheál Ó Searcóid
Lecturer: Professor Ray Bates*

Time Allowed: 2 hours

Instructions for Candidates

Answer **three (3)** questions.
All questions carry equal marks. Total: 75 marks.
Please do not use red pen on the answer books.
Please use separate answer book for each question.

Instructions for Invigilators

Non-programmable calculators may be used during this examination.

Question 1. [Marks. (a): 5. (b): 10. (c): 10]

- (a) Define potential temperature (θ). Show that when θ is constant with height (z) the lapse rate of temperature is given by

$$\Gamma_d \equiv -\frac{dT}{dz} = \frac{g}{c_p} \quad (1.1)$$

where g is the acceleration of gravity and c_p is the specific heat of dry air at constant pressure. (Assume the atmosphere is hydrostatically balanced.)

- (b) Given the thermodynamic energy equation in the form

$$c_v \frac{DT}{Dt} + p \frac{D\alpha}{Dt} = J \quad (1.2)$$

where c_v is the specific heat of dry air at constant volume, α is the specific volume and J the diabatic heating rate per unit mass, show that this equation can be written in terms of potential temperature as

$$\frac{D\theta}{Dt} = \frac{\theta}{c_p T} J \quad (1.3)$$

(Use the standard thermodynamic relationship $R = c_p - c_v$).

Hint: You may work backwards from (1.3) using the definition of θ if you wish.

- (c) Show that in the isobaric coordinate system eq. (1.2) can be written in the form

$$\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} - S_p \omega = \frac{J}{c_p} \quad (1.4)$$

where

$$S_p = \frac{\alpha}{c_p} - \frac{\partial T}{\partial p} \quad (1.5)$$

Note: $\omega \equiv Dp/Dt$.

Question 2. [Marks. (a): 5. (b): 13. (c): 7]

- a) State the assumptions used in constructing the shallow water model with a free surface of height $h_T(x,y,t)$, bottom topography $h_s(x,y)$ and fluid depth $h(x,y,t)$ on a β -plane.
- b) Starting from the primitive equations of motion on a β -plane and using the above assumptions, show that the governing equations for the model in question are

$$\frac{Du}{Dt} = fv - \frac{\partial \Phi_T}{\partial x} \quad (2.1)$$

$$\frac{Dv}{Dt} = -fu - \frac{\partial \Phi_T}{\partial y} \quad (2.2)$$

$$\frac{Dh}{Dt} = -h \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) \quad (2.3)$$

where $\Phi_T = gh_T$.

- c) Derive the linearized form of the above equations for small perturbations about a state of rest for the case where $h_s = 0$ and $\beta = 0$. Hence show the existence of 1D gravity-inertia wave solutions with phase speed

$$c = \pm \left(gH + \frac{f_0^2}{k^2} \right)^{1/2} \quad (2.4)$$

where H is the mean depth of the fluid and k is the wavenumber in x .

Question 3. [Marks. (a): 11. (b): 14]

- a) The quasi-geostrophic potential vorticity equation for the free-surface shallow water model on a β -plane ($f = f_0 + \beta y$) is given by

$$\frac{D_g}{Dt} \left[\nabla^2 \psi + f - \frac{\psi}{\lambda^2} \right] = -\frac{f_0}{\Phi_0} \frac{D_g \Phi'_s}{Dt} \quad (3.1)$$

where D_g / Dt is the material derivative following the geostrophic motion, ψ is the geostrophic streamfunction, ∇^2 is the horizontal Laplacian, $\lambda = \sqrt{\Phi_0 / f_0}$, Φ_0 is the mean geopotential of the free surface and Φ'_s is the geopotential of the orography (assumed small by comparison with Φ_0).

Using the linearized perturbation form of (3.1) for a resting basic state with $\Phi'_s = 0$, derive the phase speed for a 1D Rossby wave propagating in the x -direction in this model.

- b) Again using (3.1) but neglecting β , derive the solution for the perturbation streamfunction ψ' in the case of steady motion forced by a mean current \bar{u} blowing over orography of the form

$$\Phi'_s = \text{Re} \left[\hat{\Phi}_s \exp(ikx) \right] \quad (3.2)$$

Hint: Since the forcing has spatial dependence of the form $\exp(ikx)$, the solution may also be assumed to have spatial dependence of this form.

Question 4. [Marks. (a): 6. (b): 6. (c):13]

- (a) Show that in a hydrostatic atmosphere the internal and gravitational potential energies are proportional and that the sum of these two forms of energy (i.e., the total potential energy – TPE) can be written

$$TPE = \left(\frac{c_p}{c_v} \right) E_I \quad (4.1)$$

where c_p and c_v are the specific heats of air at constant pressure and constant volume, respectively, and E_I is the internal energy.

- (b) Show that the TPE of a unit column of atmosphere of uniform potential temperature θ extending from the surface ($p=p_0$) to the top of the atmosphere ($p=0$) is

$$TPE = \frac{c_p}{g} \frac{p_0 \theta}{\kappa + 1} \quad (4.2)$$

where $\kappa = R/c_p$ and g is the acceleration of gravity.

- (c) Consider two air masses of uniform potential temperatures θ_1 and θ_2 ($\theta_2 > \theta_1$) which are separated by a vertical partition. Each air mass occupies a horizontal area A and extends from the surface (where $p=p_0$ for each) to the top of the atmosphere. Show that the available potential energy for this system is given by

$$APE = \frac{c_p}{g} \frac{p_0}{\kappa + 1} \left(1 - \frac{1}{2^\kappa}\right) (\theta_2 - \theta_1) A \quad (4.3)$$

Question 5. [Marks. (a): 15. (b): 4. (c): 6]

- (a) Given the following governing equations for the two-layer model of baroclinic instability

$$\left(\frac{\partial}{\partial t} + U_1 \frac{\partial}{\partial x}\right) \frac{\partial^2 \psi_1'}{\partial x^2} + \beta \frac{\partial \psi_1'}{\partial x} = \frac{f_0}{\delta p} \omega_2' \quad (5.1)$$

$$\left(\frac{\partial}{\partial t} + U_3 \frac{\partial}{\partial x}\right) \frac{\partial^2 \psi_3'}{\partial x^2} + \beta \frac{\partial \psi_3'}{\partial x} = -\frac{f_0}{\delta p} \omega_2' \quad (5.2)$$

$$\left(\frac{\partial}{\partial t} + U_m \frac{\partial}{\partial x}\right) (\psi_1' - \psi_3') - U_\tau \frac{\partial}{\partial x} (\psi_1' + \psi_3') = \frac{\sigma \delta p}{f_0} \omega_2' \quad (5.3)$$

where $U_m = (U_1 + U_3)/2$ and $U_\tau = (U_1 - U_3)/2$, show that the rate of change of the sum of the kinetic and available potential energies is

$$\frac{d}{dt} (K' + P') = 4\lambda^2 U_\tau \overline{\psi_\tau \frac{\partial \psi_m}{\partial x}} \quad (5.4)$$

where $\psi_m = (\psi_1' + \psi_3')/2$, $\psi_T = (\psi_1' - \psi_3')/2$, $\lambda^2 = (f_0)^2 / [\sigma(\delta p)^2]$ and $(\overline{\quad})$

denotes an average over the wavelength of the disturbance.

- (b) Discuss the energetics of baroclinic waves qualitatively using a box diagram involving \overline{P} (the mean available potential energy), P' and K' .
- (c) Discuss the physical mechanism of baroclinic instability in terms of the slope of particle trajectories ('the wedge of instability').

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